

DE LA RECHERCHE À L'INDUSTRIE



EUROfusion



www.cea.fr

PROGRESSES IN MODELING THE RUNAWAY ELECTRON PHYSICS

(EUROFUSION, ER15-CEA-09)

***Y. Peysson**¹, G. Anastassiou², J.-F. Artaud¹, A. Budai³, J. Decker⁴,
O. Embréus⁵, O. Ficker⁸, T. Fulöp⁵, L. Hesslow⁵, K. Hizanidis², Y.
Kominis², T. Kurki-Suonio⁶, P. Lauber⁷, R. Lohner³, J. Mlynár⁸, E.
Nardon¹, G. Papp⁷, R. Paprok⁸, G. Pokol³, F. Saint-Laurent¹, C.
Reux¹, K. Sarkimaki⁶, C. Sommariva¹, A. Stahl⁵, G. Wilkie⁵, P.
Zestanakis²*

- 1) IRFM-CEA, Cadarache, France
- 2) National Technical University of Athens, Greece
- 3) BME – NTI, Budapest, Hungary
- 4) CRPP, Swiss Federal Institute of Technology, Switzerland
- 5) Chalmers University, Göteborg, Sweden
- 6) Aalto University, Finland
- 7) Max-Planck Institute for Plasma Physics, Garching, Germany
- 8) Institute for Plasma Physics, Prag, Czech Republic

<https://www2.euro-fusion.org/erwiki>

- During disruptions runaways sometimes form a beam of MeV electrons which can carry a significant fraction of the initial current. The runaway beam can become unstable and hit the wall, creating great damages → *identified as a serious issue for ITER*
- The electron distribution evolves under the combined influence of several phase-space mechanisms for RE generation (collisions, E_{\parallel} acceleration, diffusion by RF waves, synchrotron reaction force, etc) and RE transport (kinetic instabilities, MHD instabilities, RMPs, turbulence)
- During disruptions, the runaway generation can be dominated by the avalanche process : runaway dynamics is highly non-linear. In addition RE population and E_{\parallel} acceleration must be self-consistently determined
- Small variations in the balance between runaway generation and transport can lead to large differences in the resulting density of runaways and the formation of a RE beam. The influence of a seed may be also critical.

A kinetic description of runaway electron dynamics is necessary

- The Work Package for Enabling Research **ER15-CEA-09** « Kinetic modelling of runaway electron dynamics » supported by EUROfusion for studying RE dynamics started in 2014, and **will officially end in December 2017**.
(*Invaluable contribution from Pr T. Fülöp of Chalmers University*)
- The WPER has gathered the effort of about 8 European universities or research institutes
- The primary aim of the kinetic modeling effort of RE dynamics is to describe the formation of the suprathermal beam, taking into account selfconsistently of transport and non-linear effects, developing synthetic diagnostics and benchmark against experiments (*available tools → kinetic solvers: LUKE, CODE; tokamak simulators: METIS, GO, synthetic diagnostic: R5-X2,...*)
- The ultimate goal is to validate robust techniques for RE mitigations, but also find solutions to limit the formation of a multi-MeV beam once the thermal quench as occurred.

<https://www2.euro-fusion.org/erwiki> for reports and references

RE theory

Chalmers University : O. Embreus, T. Fülöp, A. Stahl, G. Wilkie, L. Hesslow

RE experiments

JET : C. Reux
TCV : S. Coda, J. Decker
COMPASS : J. Mlynar, R. Paprok
ASDEX-U: G. Papp
C-Mod: A. Tinguely, R. Granetz

RE modelling

CODE (+ GO) → **Chalmers University** : A. Stahl and **IPP Garching** : G. Papp
LUKE (+ METIS) → **CEA, France**: Y. Peysson, J.-F. Artaud
LUKE (+ ITM) → **BME, Budapest** : A. Budai, G. Pokol and **CEA, France**: Y. Peysson

RE transport

IPP Garching : G. Papp
Aalto University, Finland: T. Kurki-Suonio, K. Sarkimaki

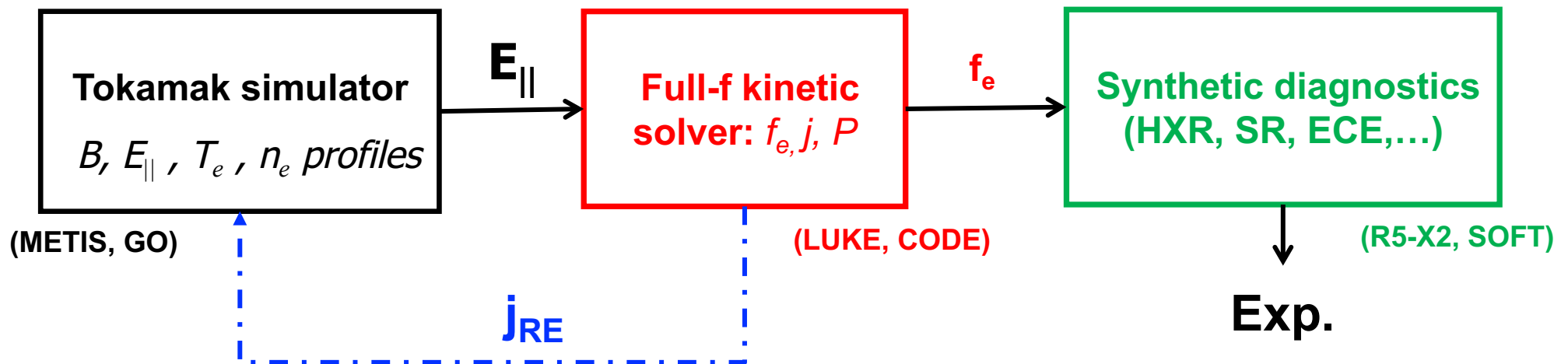
Kinetic instabilities

BME, Budapest : A. Budai, G. Pokol
NTUA, Greece: A. Zestanakis, Y. Kominis, G. Anastassiou, K. Hizanidis

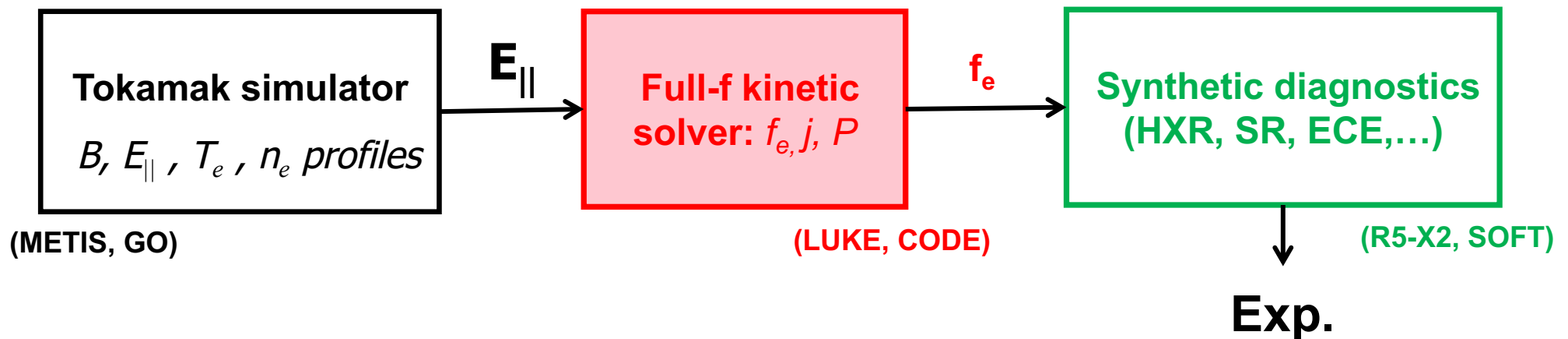
MHD instabilities

IPP Garching : G. Papp, P. Lauber
IRFM/Polytechnique : X. Garbet , R. Sabot, G. Brochard

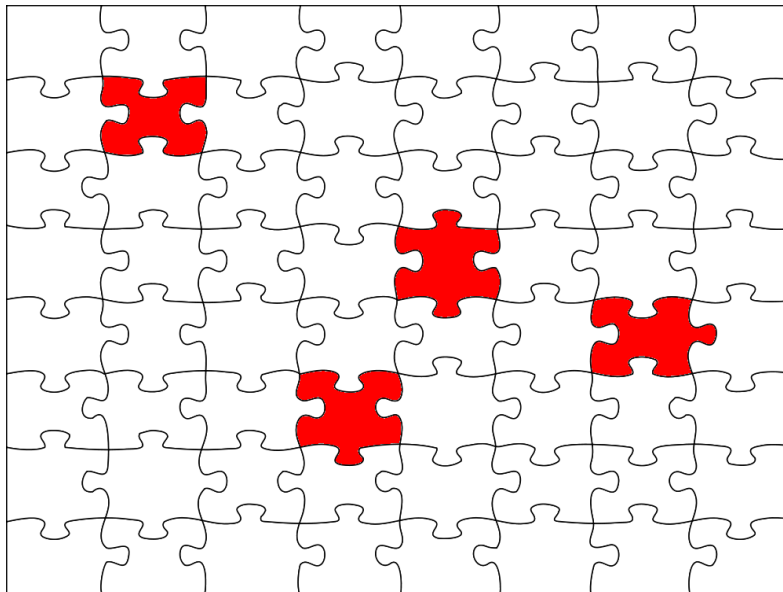
- Focusing on the generation and transport mechanisms (avalanches, additional forces that could limit the formation of a runaway beam,...)
- Integrating the various processes self-consistently in tokamak simulations
- Building synthetic diagnostics and performing comparisons with experiments.



- **Focusing on the generation and transport mechanisms (avalanches, additional forces that could limit the formation of a runaway beam,...)**
- Integrating the various processes self-consistently in tokamak simulations
- Building synthetic diagnostics and performing comparisons with experiments.



Finite difference Fokker-Planck solvers

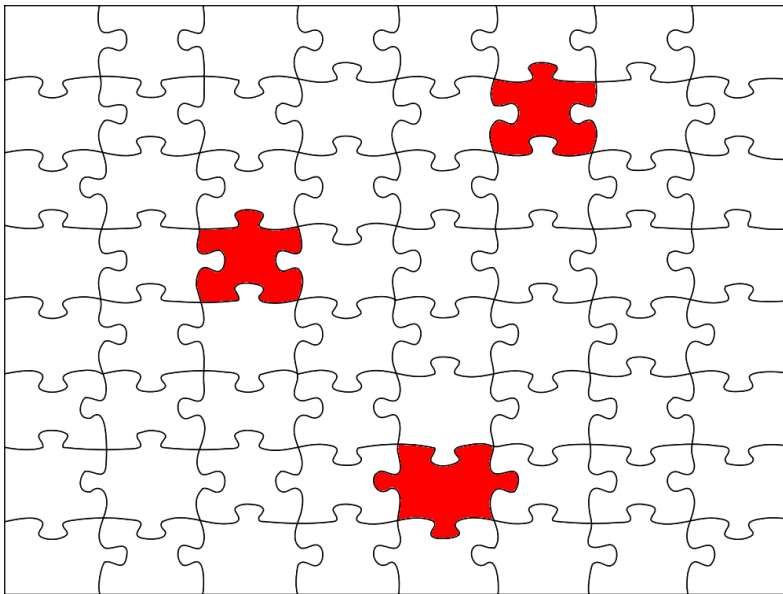


- Synchrotron radiation reaction force (ALD): CODE, LUKE (with toroidal effects)
- Bremsstrahlung radiation reaction force: CODE
- Knock-on operator for RE avalanches : CODE, LUKE (with toroidal effects)
- Effects of partially screened impurities on the collision operator : CODE
- *Implement the quasilinear model for kinetic instabilities (EXEL waves): LUKE*

CODE (cylindrical): 2-D momentum space (Chalmers U., Sweden)

LUKE (toroidal): 2-D momentum + 1D configuration spaces (CEA, France)

Finite difference Fokker-Planck solvers



- Drift-diffusion RE transport model in magnetic fields that contain both stochastic regions and islands from orbit-following code ASCOT. Applicable for LUKE in the presence of MHD instabilities, magnetic turbulence or RMP

LUKE (toroidal): 2-D momentum + 1D configuration spaces (CEA, France)

2013

$$\frac{\partial f}{\partial t} = C_{FP}(f) + Q_{RF}(f) + E(f) + T(f) + S(f)$$

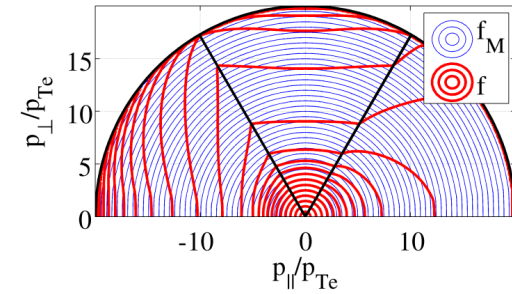
collisions → $C_{FP}(f)$ $E_{||}$ acc. → $E(f)$ Sources & sinks → $S(f)$
 RF heating → $Q_{RF}(f)$ radial trans. → $T(f)$

2017

$$\frac{\partial f}{\partial t} = C_{FP}(f) + C_{KO}(f) + Q_{RF}(f) + E(f) + R(f) + T(f) + S(f)$$

collisions → $C_{FP}(f)$ and $C_{KO}(f)$ $E_{||}$ acc. → $E(f)$ Sources & sinks → $S(f)$
 RF heating → $Q_{RF}(f)$ synchrotron reac. → $R(f)$ radial transp. → $T(f)$

- **Full-f finite difference code with non-uniform grids**
 - Energy scale from eV to hundred MeV
 - Refined description near $\xi_0 = p_{||0}/p \sim 1$
 - Radial zoom over regions of interest
- **Various geometries** (toroidal, dipole, cylindrical) ; plasma shaping effect. Zero banana width approximation (ZOW)
- Choice of **different time schemes**: from collisional to equilibrium times
- **Standard moments of f_e** : J , P_{RF} , P_C , P_E , Γ_R , ...
- **Benchmarked** against neoclassical conductivity, RF heating, primary and secondary runaway rate, etc. Extensively used for LH, EC and EBW physics
- **Numerical structure ready to include self-consistent radial and cross-term dynamics (FOW)**: neoclassical transport, wave-induced transport, Ware pinch, etc
- Self-consistent ripple losses: loss cone and critical energy



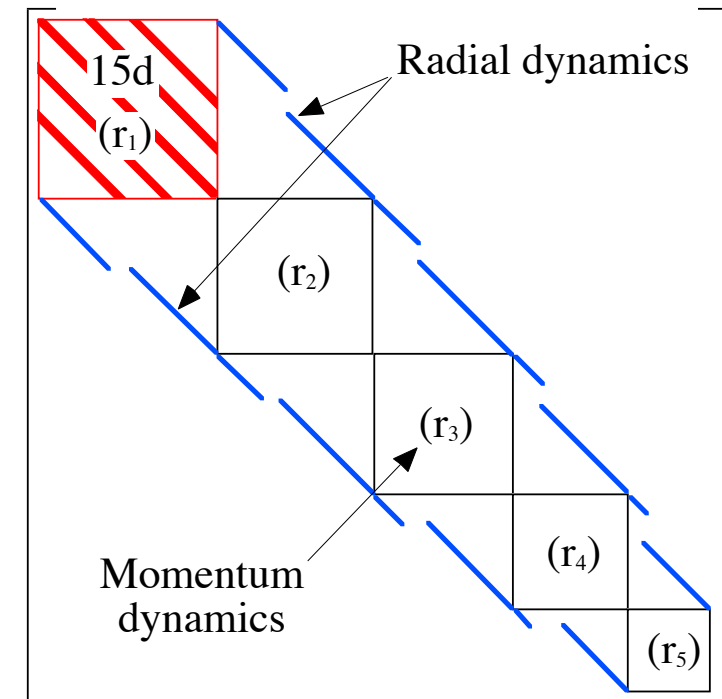
- **Bounce-averaged Fokker-Planck equation is cast in a conservative form**

$$\frac{\partial f}{\partial t} + \nabla \cdot \mathbf{S}_{\psi, p, \xi_0}(f) = K(f; \psi, p, \xi_0)$$

- **Fully relativistic Fokker-Planck collision operator including momentum conserving integral term**

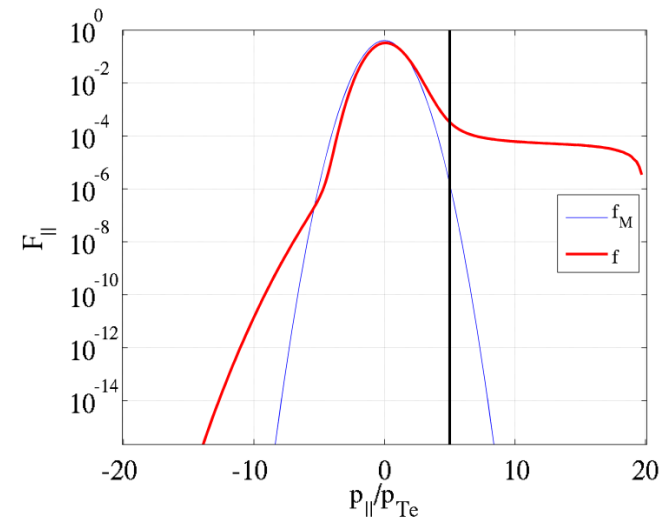
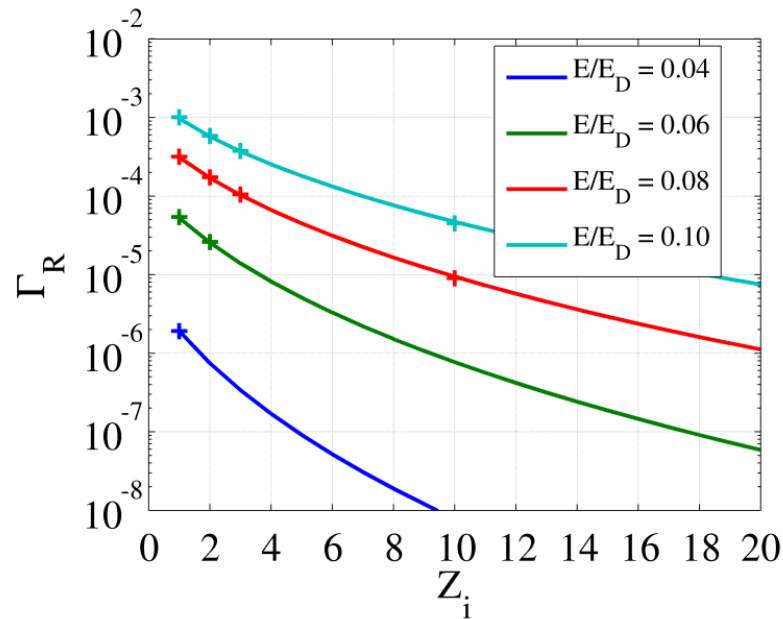
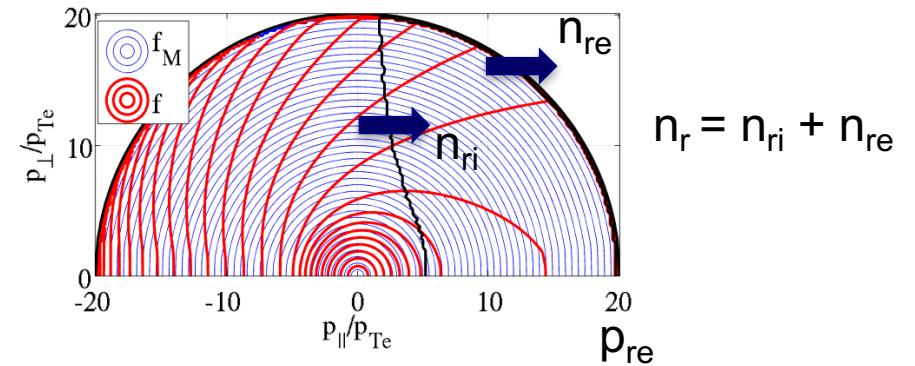
$$C(f) = \nabla \cdot \mathbf{S}^C(f) + I^C(f)$$

- **Universal quasilinear operator for electron interaction with all types of RF waves**
- Written in Matlab® and coupled to powerful external matrix inversion packages (MUMPS, ...) with some specific modules in C (C3PO ray tracing) or Fortran.
- Distributed (CPU+ GPU) and remote processing for fast calculations



Peysson, Y. and Decker, J., FST, 65 (2014) 22

- Standard runaway rate calculation (Dreicer mechanism)
- Non-relativistic limit, cylindrical plasma
- Excellent agreement with Kulsrud theory (+)



R.M. Kulsrud et al., Phys. Rev. Lett., 1973, 31, 11, pp. 690-693

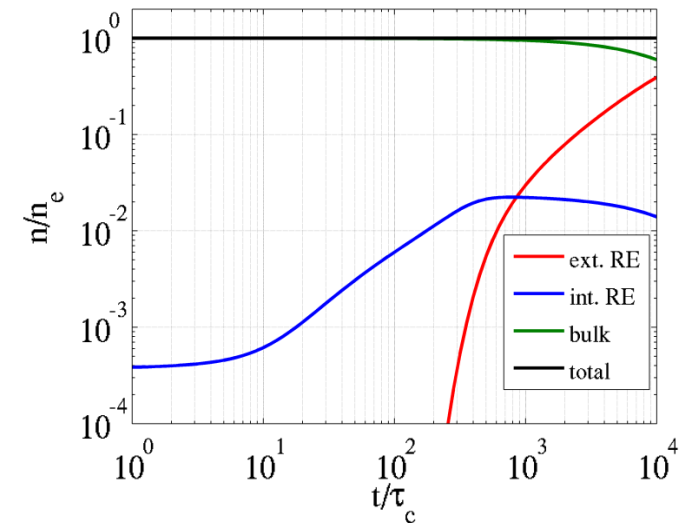
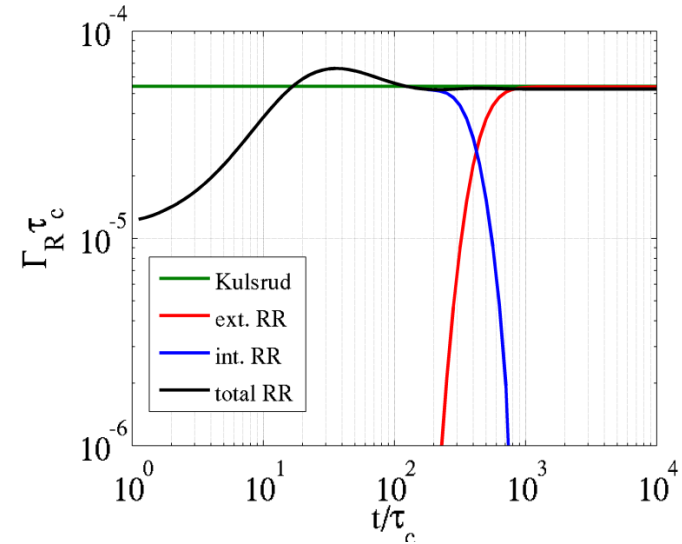
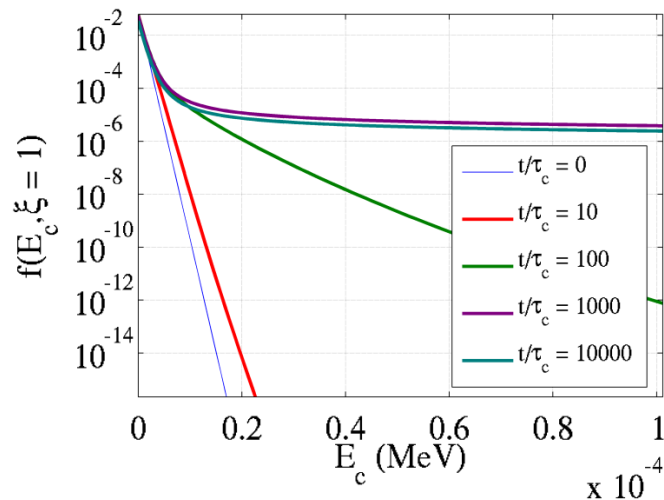
- $E/E_D = 0.06$, $Z_i = 1$, NR limit, cylindrical plasma

- Primary runaway rate :

$$\Gamma_r = \frac{1}{n_b} \frac{dn_r}{dt} = \frac{1}{n_b} \frac{dn_{ri}}{dt} + \frac{1}{n_b} \frac{dn_{re}}{dt}$$

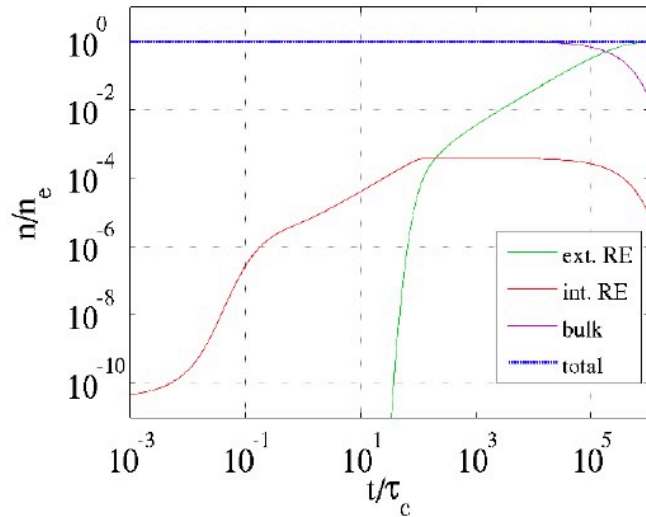
- Time asymptotic rate :

$$\Gamma_r = \frac{1}{(1 - n_{re})} \frac{dn_{re}}{dt}$$



LUKE, CODE : time evolution of primary runaway rate with synchrotron radiation reaction force (ALD)

NO ALD FORCE



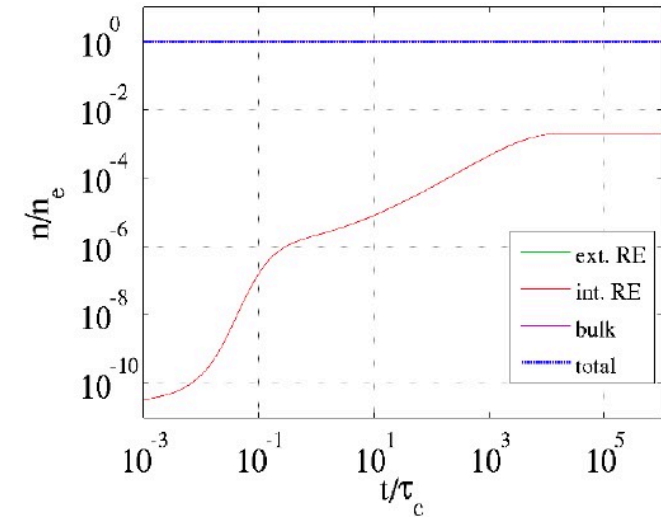
$$E_{\parallel} = 3$$

$$\sigma_r = 0.6$$

$$\beta = 0.1$$

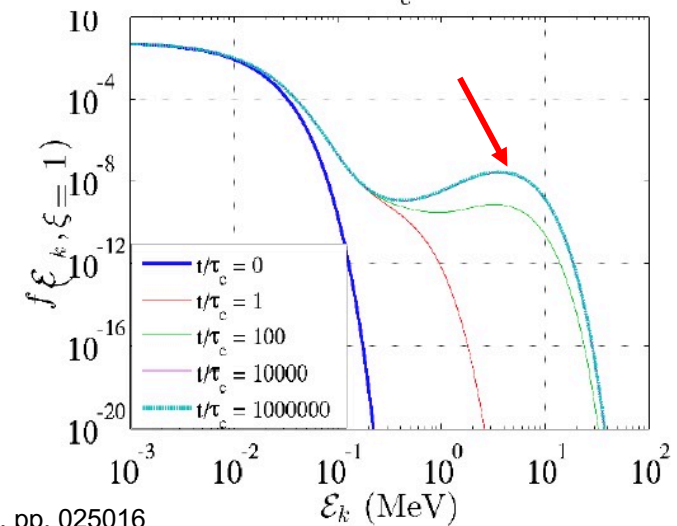
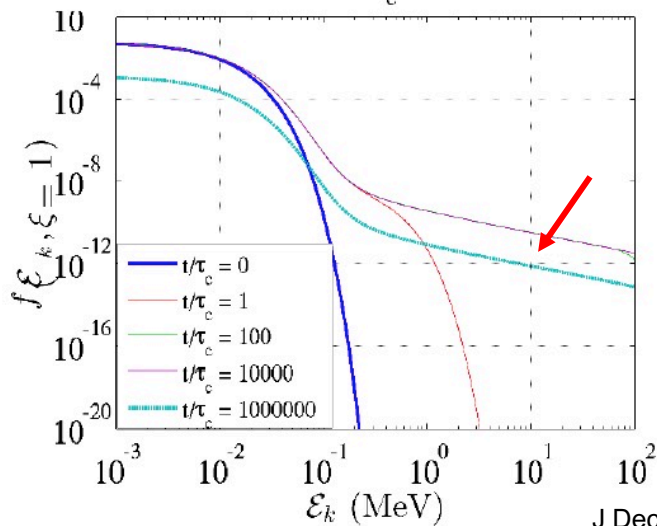
$$Z_i = 1$$

WITH ALD FORCE



ALD

- Upper energy limit of ten MeV
- Steady state regime



J Decker et al., *Plasma Phys. Control. Fusion*, 2016, 58, pp. 025016

- Rosenbluth-Putvinski formulation using Moller relativistic electron-electron differential cross section
→ **Maxwellian electrons are at rest.** Knock-on electrons described by a source term in FP codes

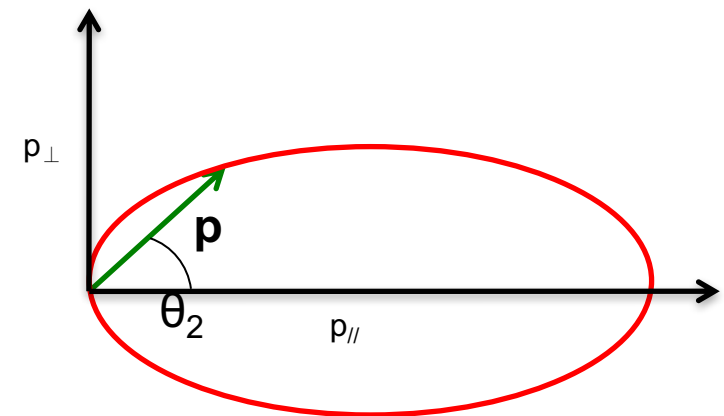
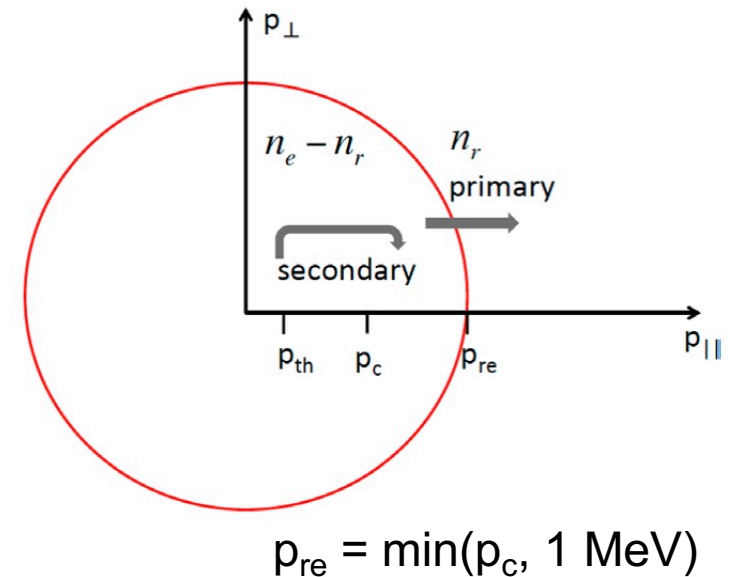
$$S(p, \psi, \xi, \theta) = n_e n_R c \frac{d\sigma}{d\Omega} = n_e n_R c r_e^2 \frac{1}{p\gamma(\gamma - 1)^2} \delta(\xi - \xi^*(p))$$

- The knock-on electrons emerge highly magnetized
→ **bounce-averaging is essential**

$$\{\bar{S}\} = \frac{1}{\lambda \tilde{q}} \left[\frac{1}{2} \sum_{\sigma} \right]_T \int_{\theta_{min}}^{\theta_{max}} \frac{d\theta}{2\pi} \frac{1}{|\hat{\psi} \cdot \hat{r}|} \frac{r}{R_p} \frac{B}{B_p} \frac{\xi_0}{\xi} \bar{S}(\psi, p, \xi)$$

- Particle conserving form of avalanche process:

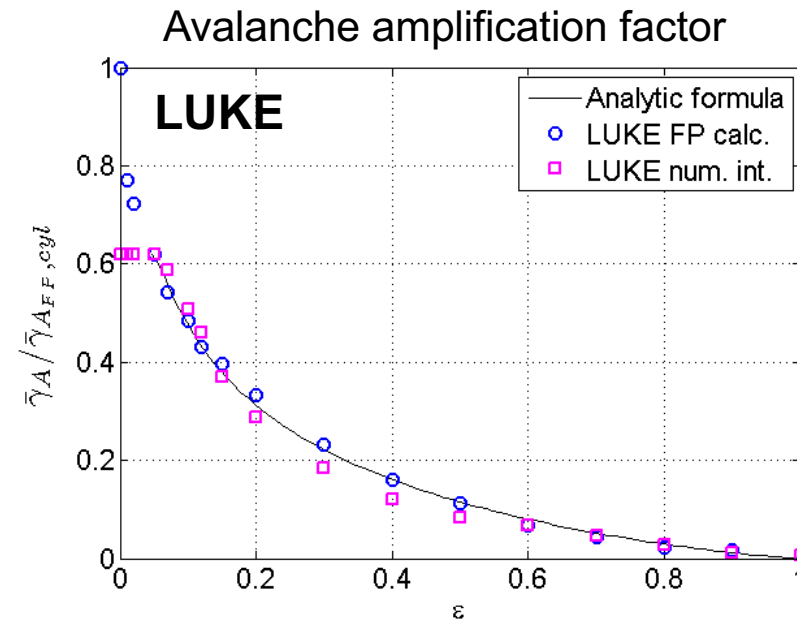
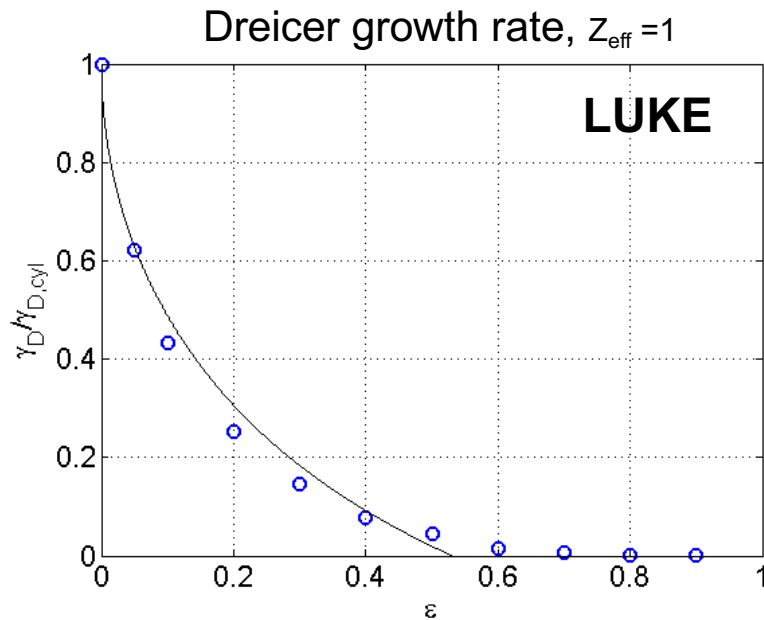
$$S = S_+ - \langle S_+ \rangle \bar{f}_M \quad \left\{ \begin{array}{l} S_+ = n_e n_r c \frac{d\sigma}{d\Omega} \\ n_r + n_e = const \end{array} \right.$$



E.Nilsson et al., Plasma Phys. Control. Fusion, 2015, 57, 9, pp. 095006

Growth rate: $\frac{\partial n_r}{\partial t} = n_e (\gamma_D + \gamma_A)$

$\gamma_A = n_r \bar{\gamma}_A$



E.Nilsson et al., Plasma Phys. Control. Fusion, 2015, 57, 9, pp. 095006

Runaway rate strongly reduced due to trapped electrons

Agrees with predictions by ARENA code [Eriksson & Helander, *Comp. Phys. Comm.* **154** (2003)]

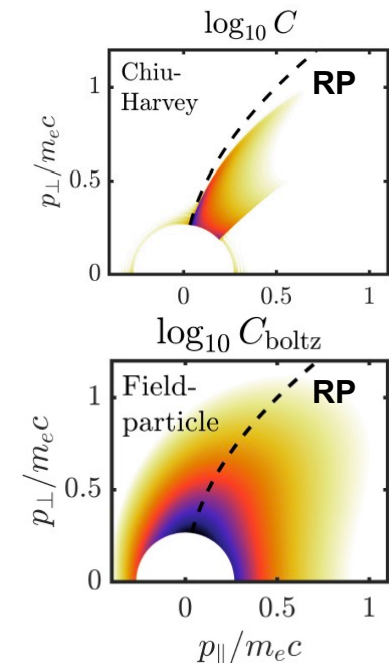
and CQL3D code [Harvey & McCoy, *IAEA* (1992)] [Decker & Peysson, *EUR-CEA-FC-1736, Euratom-CEA*, 2004]

- To describe knock-on collisions a (simplified) linearized Boltzmann operator is added ($n_r \ll C_{\text{knock-on}} = C_{\text{boltz}} \{n_e \delta(\mathbf{p}), f_e\}$ (only field-particle term))

Rosenbluth-Putvinski: $f_e(\mathbf{p}) = n_{\text{RE}} \lim_{p_0 \rightarrow \infty} \frac{1}{p^2} \delta(p - p_0) \delta(\cos \theta - 1)$

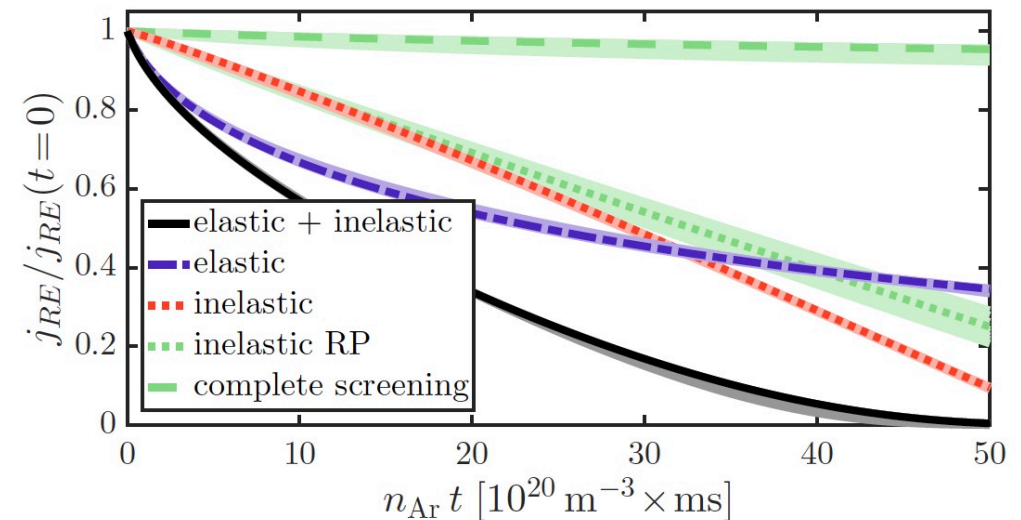
Chiu-Harvey: $f_e(\mathbf{p}) = F(p) \delta(\cos \theta - 1)$ $F(p) = \int_{-1}^1 f_e(\mathbf{p}) d(\cos \theta)$

- Model limitation:
 - Double counting collisions
 - Non-conservation of momentum and energy
 - Chiu-Harvey model ignores pitch-angle distribution
 - Arbitrary cut-off affecting solutions
- New rigorous approach:**
 - Accounting for full $f_e(p)$
 - Including the test-particle term: *restores conservation laws*
 - Modify $\ln \Lambda$ in Fokker-Planck operator: *avoids double counting*



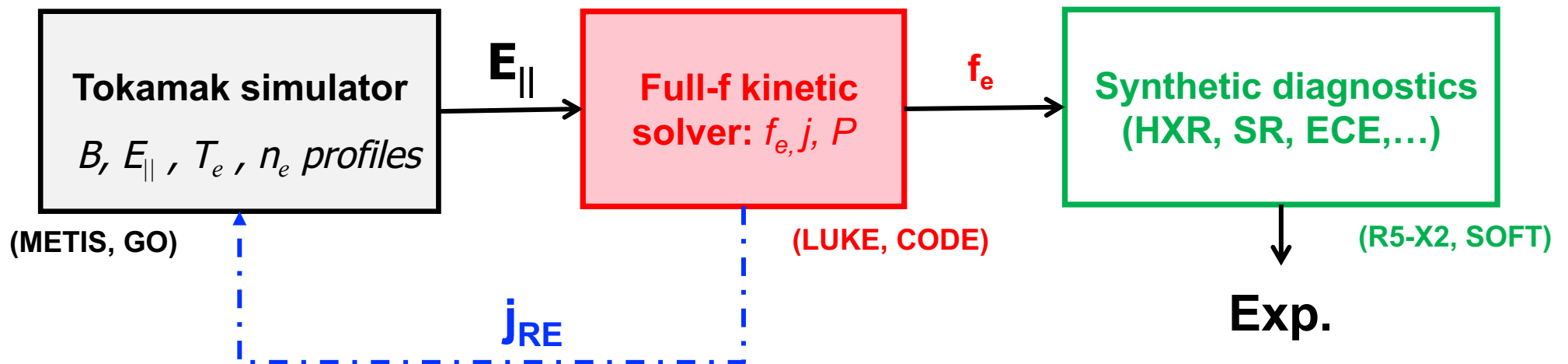
Reduction of the avalanche generation for $E/E_c \leq 2-3$

- Weakly ionized impurities with $Z \gg 1$ in very cold plasma (Ar: 18, W: 74)
- Elastic scattering e/i described in Born approximation, $v/c \gg Z\alpha$ ($\alpha = 1/137$)
- Weakly energetic electrons “feels” a screened ion, while very energetic ones interacts with the nucleus → more accurate collision operator
- Inelastic scattering e/e needs also corrections
- **Large effect of reduced screening (CODE)**
 - Enhanced collision frequencies
 - More isotropic distribution function
 - Enhanced effective critical electric field E_c^*
 - Faster runaway current decay

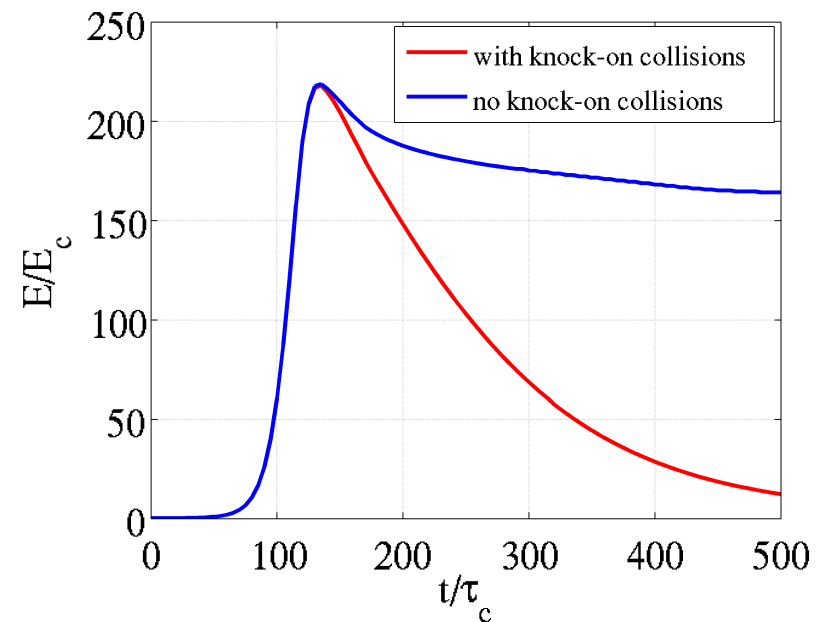
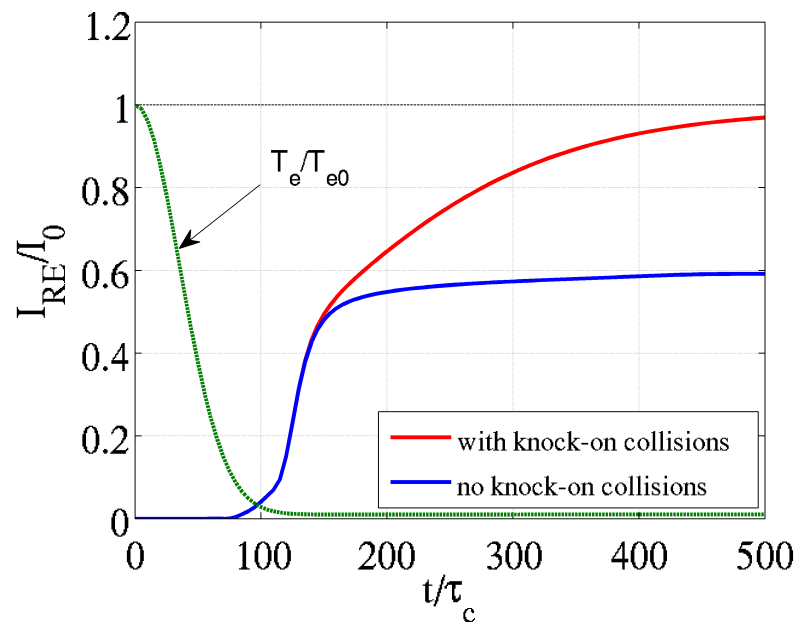


L. Hesslow et al., Phys. Rev. Letter, 2017, 118, pp. 255001

- Focusing on the generation and transport mechanisms (avalanches, additional forces that could limit the formation of a runaway beam,...)
- Integrating the various processes self-consistently in tokamak simulations
- Building synthetic diagnostics and performing comparisons with experiments.

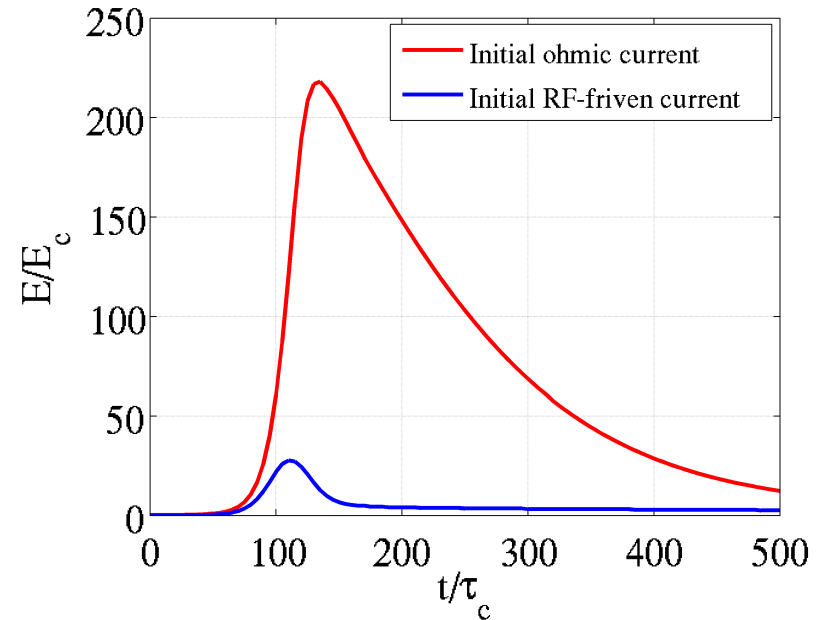
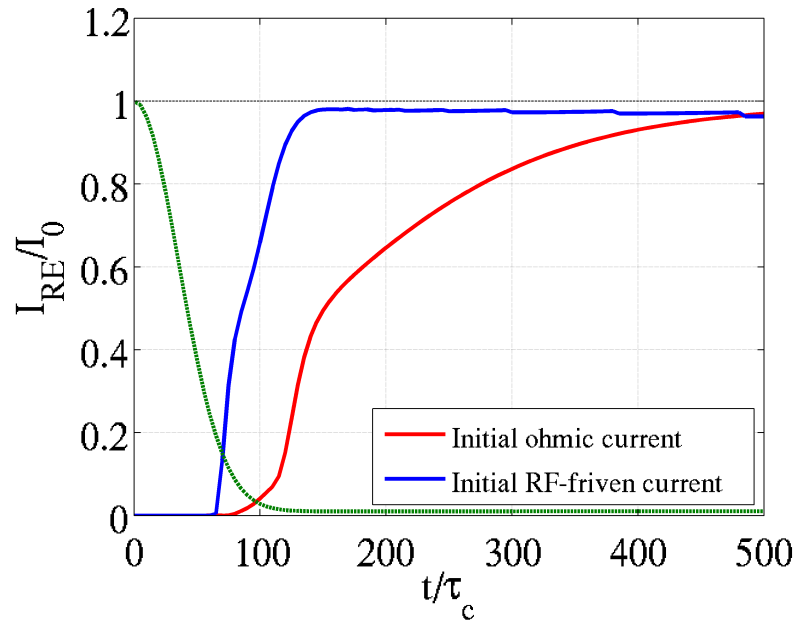


- $Z_i = 1$, cylindrical plasma, **local electric field adjusted to constant $J=J_0 = 1 \text{ MA/m}^2$**
- **T_e evolution enforced through collisions: $T_e(t) = T_f + (T_{e0} - T_{ef}) \times \exp[-(t/t_q)^2]$**
- $T_{e0} = 5.1 \text{ keV}$; $T_{ef} = 51 \text{ eV}$; $t_q = 50 \tau_c$ (5 ms)

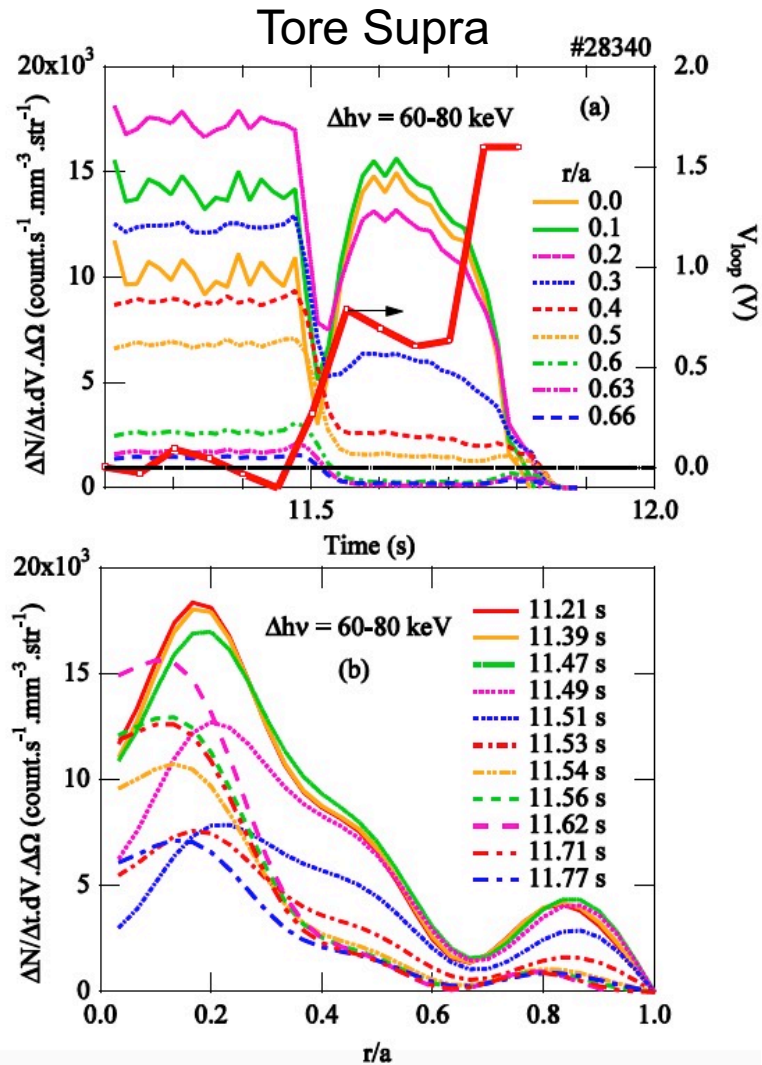


- Avalanche effect dominates E-field evolution (R-P model not so bad as $E/E_c \gg 1$)
- Effect of synchrotron reaction force is negligible (internal RE, not energetic enough)

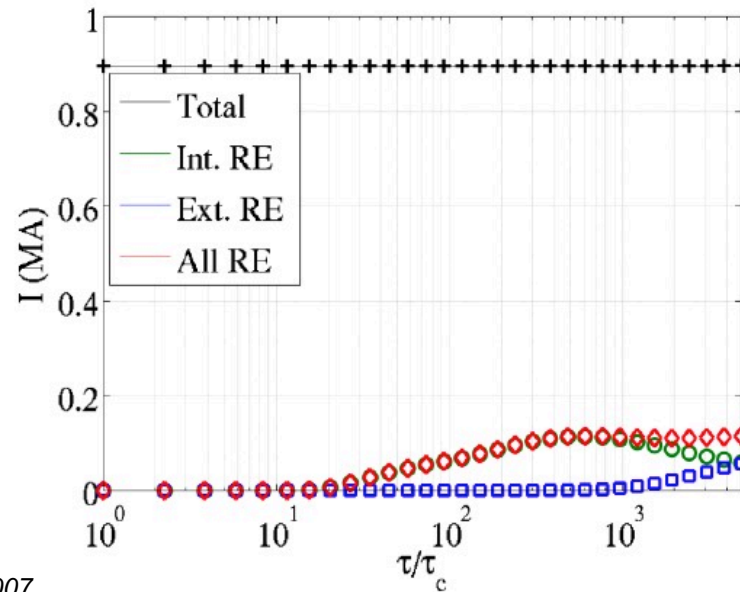
A. Stahl et al., Nucl. Fusion, 2016, 56, pp. 112009



- LHCD, $v_{\min} = 3.5$, $v_{\max} = 7.0$, P_{LH} adjusted to $J_0 = 1 \text{ MA/m}^2$, LH power turned off at $t = 0$
- With LH-driven initial distribution, knock-on collisions play no significant role (but R-P avalanche model likely inaccurate since $E/E_c \sim 1$)
- Importance of seed fast electrons



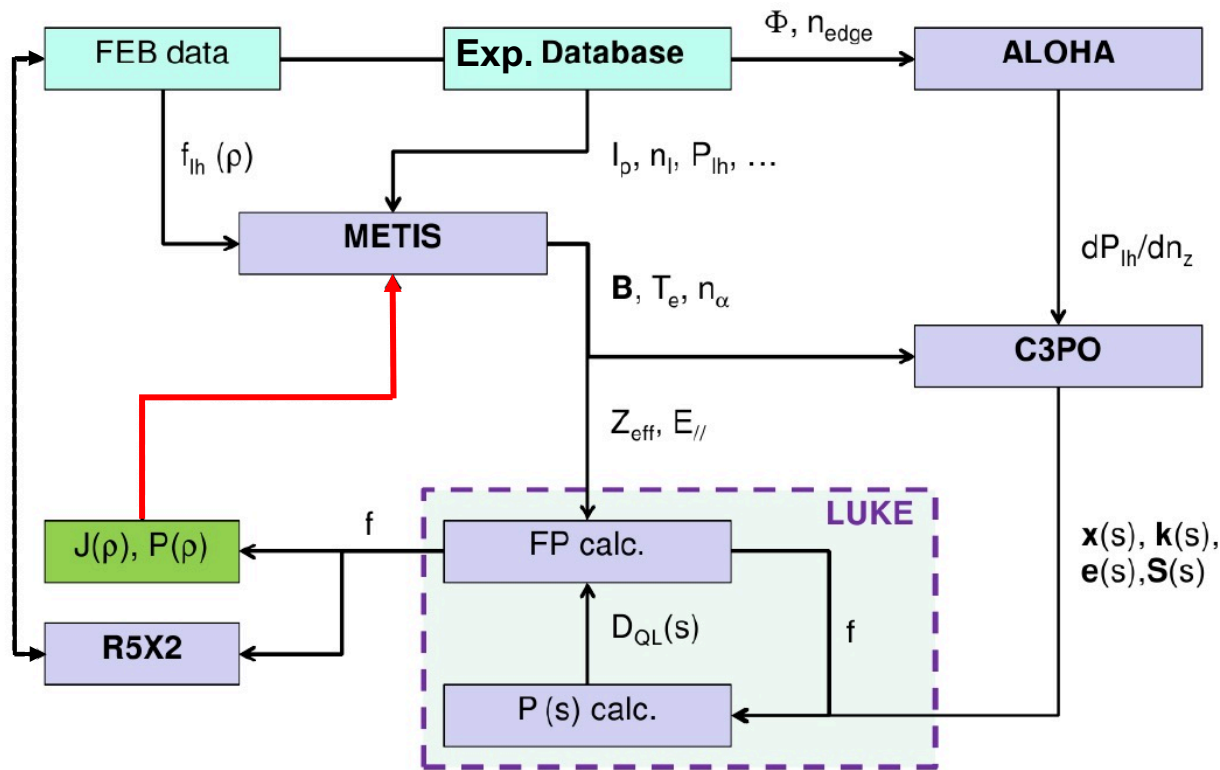
- Sudden drop of LH power from 4.3 MW (steady-state non-inductive discharge) to 1.0 MW then 0 MW at constant $I_p = 0.85$ MA, $B_t = 3.7$ T
- Time-space resolved HXR bremsstrahlung



18% of current is carried by RE

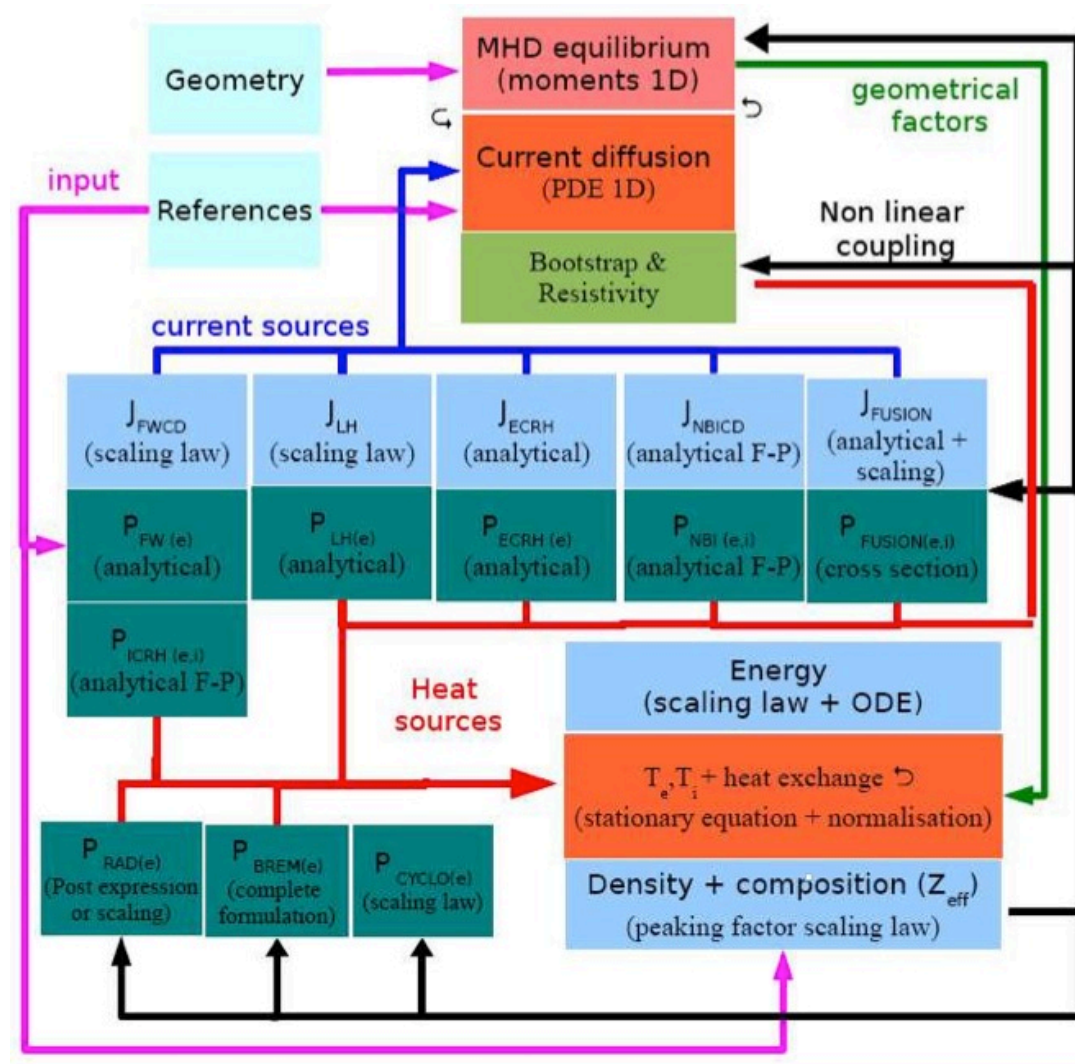
Y. Peysson and F. Imbeaux, *Rev. Sci. Instr.*, 1999, 70, 10, pp. 3987-4007

- Cylindrical : GO (1D radial) + CODE (2D momentum space)
- **Toroidal, arbitrary plasma shape : METIS (1D radial) + LUKE (1D radial + 2D momentum space)**
- LUKE in ITM and soon in **IMAS (ITER)**



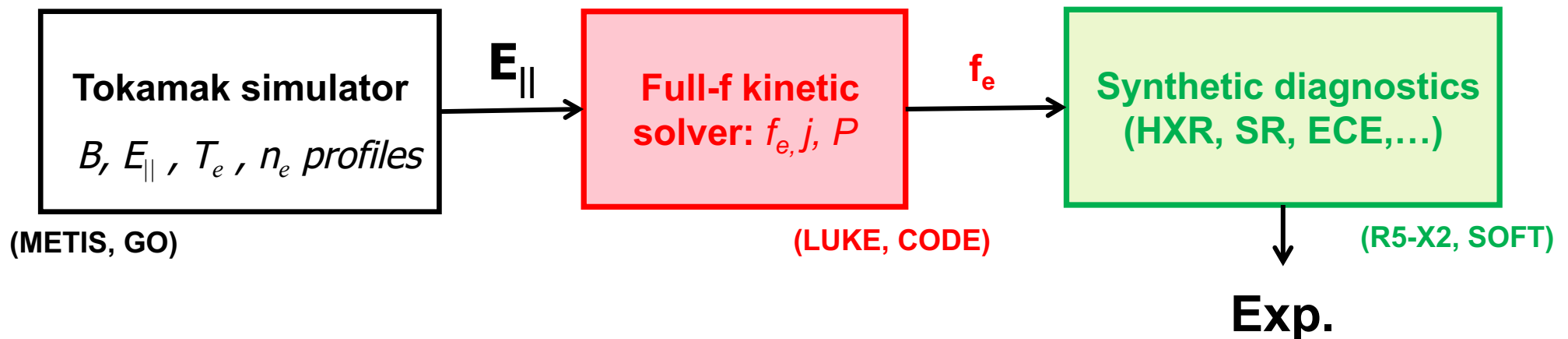
- **Full coupling loop METIS – LUKE is done**
- Stability of the convergence between j and $E_{//}$ is under investigation.
- Test Tore Supra discharge #28340
- *Effect of RF waves on RE can be studied (EC waves)*

G. Papp et al., Nucl. Fusion, 2013, 53, pp. 123017
 Artaud et al., Nucl. Fusion, 2010, 50, 4, pp. 043001



Matlab
Simulink

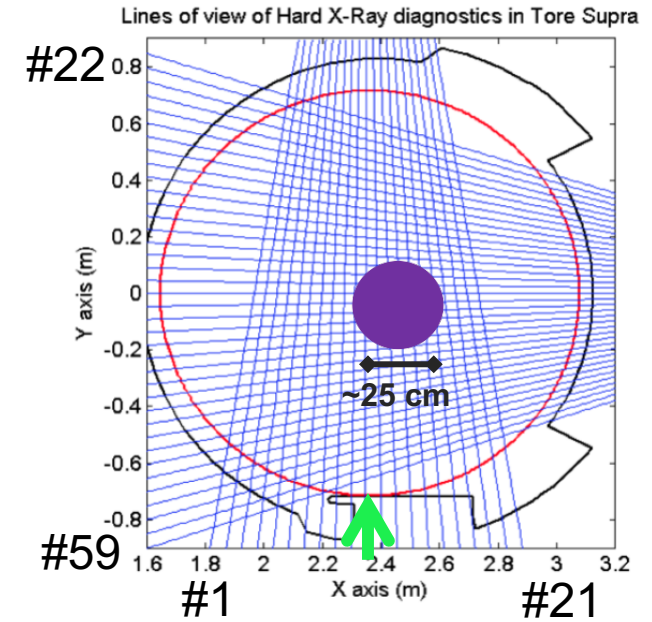
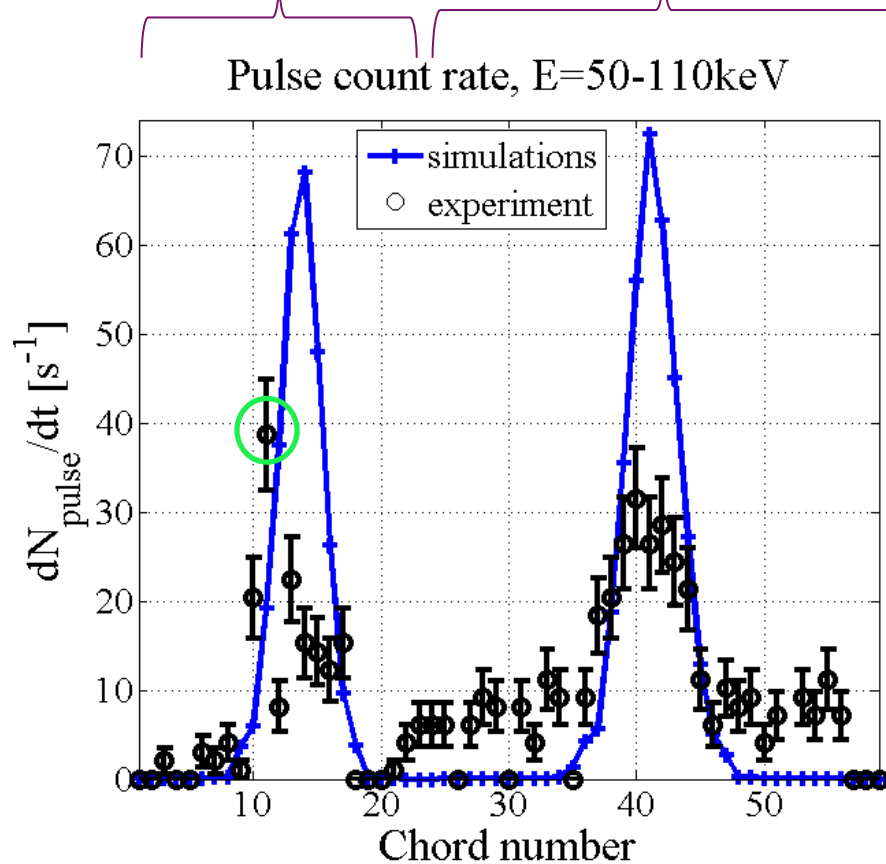
- Focusing on the generation and transport mechanisms (avalanches, additional forces that could limit the formation of a runaway beam,...)
- Integrating the various processes self-consistently in tokamak simulations
- Building synthetic diagnostics and performing comparisons with experiments.



Fast electron bremsstrahlung emission from RE (Tore Supra)

- HXR tomographic system
- **R5X2**: Synthetic diagnostic for bremsstrahlung emission [Peysson & Decker, *Phys. of Plasmas* **15** (2008)]

Vertical cam, 1-21 Horizontal cam, 22-59



RE discharge: Emission profile reproduced at the end of the current flattop (METIS+LUKE (open loop) + R5-X2)

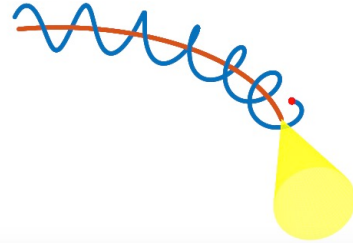
← X-rays backscattered by the tokamak inner wall [Peysson et al., *Nucl. Fusion* **33** (1993)]

E. Nilsson, PhD thesis, Ecole Polytechnique, France, 2015

- New synthetic diagnostic **SOFT**
- Angular and spectral distribution of synchrotron radiation:

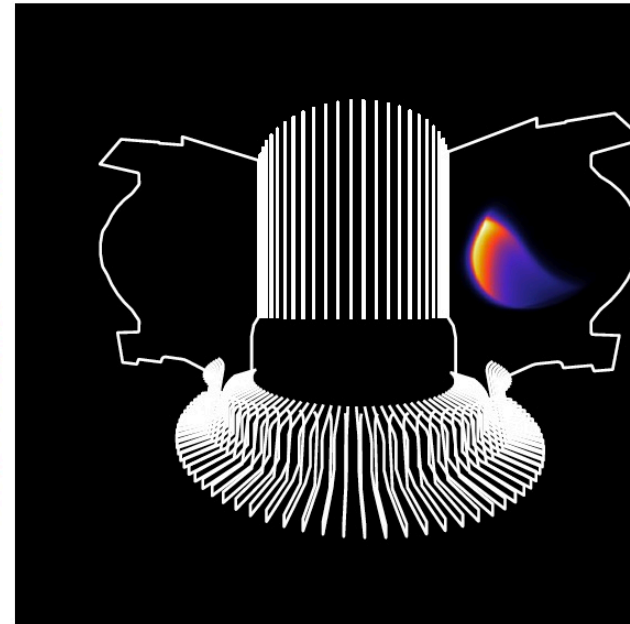
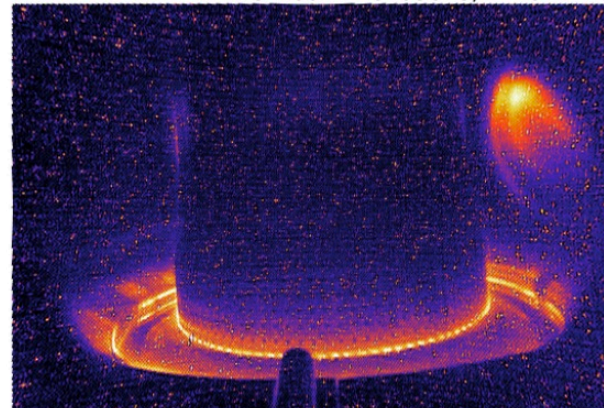
$$\frac{d^2P}{d\omega d\Omega} = \frac{3e^2\beta^2\gamma^6\omega_B}{32\pi^3\epsilon_0 c} \left(\frac{\omega}{\omega_c}\right)^2 \left(\frac{1-\beta\cos\psi}{\beta\cos\psi}\right)^2 \times \left[K_{2/3}^2(\xi) + \frac{(\beta/2)\cos\psi\sin^2\psi}{1-\beta\cos\psi} K_{1/3}^2(\xi) \right]$$

- Gyro-angle averaging
- Alcator C-Mod, 3-8 T



C-Mod 1140403026, t ~ 0.742 s

$E_{RE} \approx 30 \text{ MeV}$



- Within the Work Package for Enabling Research framework, significant improvements have been achieved in four years for realistic post-disruptive RE kinetic modeling. Progresses have been essentially devoted to momentum space physics (with and without toroidal corrections). A synthetic diagnostic has been developed for synchrotron radiation (SOFT).
- Full 1-D effects taking into account of toroidal curvature and plasma shaping is the next challenge of RE kinetic modeling: transport in mixed stochastic/coherent structures (link with MHD and orbit-following codes), as well as interaction between RE and MHD instabilities. Validation of the modeling tools against experiments must be also performed at large scale with the developed tools.
- The LUKE Fokker-Planck solver coupled to METIS simulator is a powerful tool for purpose. *Contributions from non-CEA collaborators within CEA license are welcome (use and developments)*
- Even if the Enabling Research framework will stop by end of 2017 (except if a new call from EUROfusion is launched very soon!), the activity on RE physics modeling will continue through active bilateral collaborations. A another success of the WPENR !